

Factored Lorentz Matrix Via Infeld-Van Der Waerden Symbols

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Abstract: We use the Infeld-van der Waerden symbols to factorize an arbitrary Lorentz matrix.

Key words: Lorentz matrix - Infeld-van der Waerden symbols - Unimodular matrix - Special relativity

INTRODUCTION

We know that any Lorentz matrix can be generated with the Infeld-van der Waerden symbols [1-4] via the following expression [5-8]:

$$L^\mu{}_\nu = \sigma^\mu{}_{A\dot{R}} B^A{}_C \sigma_\nu{}^{C\dot{E}} B_{\dot{E}}^{\dagger R}, \quad B = \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \end{pmatrix}, \quad B^\dagger = \begin{pmatrix} \alpha^* & \gamma^* \\ \beta^* & \delta^* \end{pmatrix}, \quad \alpha \delta - \beta \gamma = 1, \quad (1)$$

Then:

$$L^\mu{}_\nu = \sigma^\mu{}_{A1} B^A{}_C \sigma_\nu{}^{C\dot{E}} B_{\dot{E}}^{\dagger 1} + \sigma^\mu{}_{A2} B^A{}_C \sigma_\nu{}^{C\dot{E}} B_{\dot{E}}^{\dagger 2} = {}_1P^\mu{}_C {}_1Q^C{}_\nu + {}_2P^\mu{}_C {}_2Q^C{}_\nu, \quad (2)$$

with the rectangular matrices:

$$({}_rP^\mu{}_C)_{4 \times 2} = (\sigma^\mu{}_{A\dot{r}} B^A{}_C), \quad ({}_rQ^C{}_\nu)_{2 \times 4} = (\sigma_\nu{}^{C\dot{E}} B_{\dot{E}}^{\dagger r}), \quad r = 1, 2. \quad (3)$$

It is simple to construct the matrices (3), in fact:

$$({}_1P^\mu{}_C) = \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha & \beta \\ \gamma & \delta \\ i\gamma & i\delta \\ \alpha & \beta \end{pmatrix}, \quad ({}_1Q^C{}_\nu) = \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha^* & \beta^* & i\beta^* & \alpha^* \\ \beta^* & \alpha^* & -i\alpha^* & -\beta^* \end{pmatrix}, \quad (4)$$

$$({}_2P^\mu{}_C) = \frac{1}{\sqrt{2}} \begin{pmatrix} \gamma & \delta \\ \alpha & \beta \\ -i\alpha & -i\beta \\ -\gamma & -\delta \end{pmatrix}, \quad ({}_2Q^C{}_\nu) = \frac{1}{\sqrt{2}} \begin{pmatrix} \gamma^* & \delta^* & i\delta^* & \gamma^* \\ \delta^* & \gamma^* & -i\gamma^* & -\delta^* \end{pmatrix},$$

Hence from (2) and (4) the Lorentz matrix is factored in the form:

$$(L^\mu{}_\nu) = ({}_1P \quad {}_2P) \begin{pmatrix} {}_1Q \\ {}_2Q \end{pmatrix} = \frac{1}{2} \begin{pmatrix} \alpha & \beta & \gamma & \delta \\ \gamma & \delta & \alpha & \beta \\ i\gamma & i\delta & -i\alpha & -i\beta \\ \alpha & \beta & -\gamma & -\delta \end{pmatrix} \begin{pmatrix} \bar{\alpha} & \bar{\beta} & i\bar{\beta} & \bar{\alpha} \\ \bar{\beta} & \bar{\alpha} & -i\bar{\alpha} & -\bar{\beta} \\ \bar{\gamma} & \bar{\delta} & i\bar{\delta} & \bar{\gamma} \\ \bar{\delta} & \bar{\gamma} & -i\bar{\gamma} & -\bar{\delta} \end{pmatrix}, \quad (5)$$

In accordance with [9-13]; therefore [7, 14-19]:

$$\begin{aligned}
 L^0_0 &= \frac{1}{2}(\alpha\bar{\alpha} + \beta\bar{\beta} + \gamma\bar{\gamma} + \delta\bar{\delta}), & L^1_0 &= \frac{1}{2}(\bar{\alpha}\gamma + \bar{\beta}\delta) + cc, & L^2_0 &= -\frac{i}{2}(\alpha\bar{\gamma} - \bar{\beta}\delta) + cc, \\
 L^0_1 &= \frac{1}{2}(\bar{\alpha}\beta + \bar{\gamma}\delta) + cc, & L^1_1 &= \frac{1}{2}(\bar{\alpha}\delta + \beta\bar{\gamma}) + cc, & L^2_1 &= -\frac{i}{2}(\alpha\bar{\delta} + \beta\bar{\gamma}) + cc, \\
 L^0_2 &= -\frac{i}{2}(\bar{\alpha}\beta + \bar{\gamma}\delta) + cc, & L^1_2 &= -\frac{i}{2}(\bar{\alpha}\delta + \beta\bar{\gamma}) + cc, & L^2_2 &= \frac{1}{2}(\bar{\alpha}\delta - \bar{\beta}\gamma) + cc, \\
 L^0_3 &= \frac{1}{2}(\alpha\bar{\alpha} - \beta\bar{\beta} + \gamma\bar{\gamma} - \delta\bar{\delta}), & L^1_3 &= \frac{1}{2}(\bar{\alpha}\gamma - \bar{\beta}\delta) + cc, & L^2_3 &= -\frac{i}{2}(\alpha\bar{\gamma} + \bar{\beta}\delta) + cc, \\
 L^3_0 &= \frac{1}{2}(\alpha\bar{\alpha} + \beta\bar{\beta} - \gamma\bar{\gamma} - \delta\bar{\delta}), & L^3_1 &= \frac{1}{2}(\bar{\alpha}\beta - \bar{\gamma}\delta) + cc, & L^3_2 &= -\frac{i}{2}(\bar{\alpha}\beta - \bar{\gamma}\delta) + cc, \\
 L^3_3 &= \frac{1}{2}(\alpha\bar{\alpha} - \beta\bar{\beta} - \gamma\bar{\gamma} + \delta\bar{\delta}), & \alpha\delta - \beta\gamma &= 1,
 \end{aligned} \tag{6}$$

where cc means the complex conjugate of all the previous terms.

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