

## Local and Isometric Embedding of $R_4$ Into $E_5$

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**Abstract:** We establish conditions whose fulfillment implies that the spacetime admits embedding into  $E_5$ .

**Key words:** Local and isometric embedding •  $R_4$  embedded into  $E_5$  • Gauss-Codazzi equations • Thompson-Schrank theorem • Churchill-Plebański classification

### INTRODUCTION

$R_4$  accepts local and isometric embedding into  $E_5$  if we can find the second fundamental form  $b_{\mu\nu} = b_{\nu\mu}$  satisfying the equations [1-6]:

$$R_{\mu\nu\alpha\beta} = \varepsilon(b_{\mu\alpha} b_{\nu\beta} - b_{\mu\beta} b_{\nu\alpha}) \text{ Gauss} \quad (1.a)$$

$$b_{\mu\nu;\alpha} = b_{\mu\alpha;\nu} \text{ Codazzi} \quad (1.b)$$

where  $\varepsilon = \pm 1$ ,  $R_{\mu\nu\alpha\beta}$  is the curvature tensor and  $\alpha$  indicates a covariant derivative.

In Sec. 2 we indicate intrinsic sufficient conditions which must be satisfied by the geometry of  $R_4$  to guarantee the existence of  $b_{\mu\nu}$ . In Sec. 3 we use the Thompson-Schrank's theorem to show that  $b_{\mu\nu} b^{\mu\nu} \neq 0$  for the second fundamental form of  $R_4$  embedded into  $E_5$ .

**Spacetimes of Class One:** We know that the Gauss equation implies the identity [7-9]:

$$p b_{\mu\nu} = \frac{K_2}{48} g_{\mu\nu} - \frac{1}{2} R_{\mu\alpha\beta\nu} G^{\alpha\beta}, \quad p = \frac{\varepsilon}{3} b_{\alpha\beta} G^{\alpha\beta}, \quad (2)$$

where  $G_{\alpha\beta}$  is the Einstein tensor and  $K_2$  is the Lanczos scalar [9-12]:

$$K_2 \equiv {}^*R_{\mu\nu\alpha\beta}^* R^{\mu\nu\alpha\beta}, \quad (3)$$

such that [13]:

$$\det(b^\mu{}_\nu) = -\frac{1}{24} K_2. \quad (4)$$

On the other hand, we have the Thomas theorem [14]:

“If  $\det(b^\mu{}_\nu)$  then (1.a) implies (1.b)”, (5)

which means that if  $R_4$  has  $K_2 \neq 0$  and we find  $b_{\mu\nu}$  satisfying the Gauss equation, then the Codazzi condition is valid.

From (2) is immediate the relation:

$$p^2 = -\frac{\varepsilon}{6} \left( \frac{R}{24} K_2 + R_{\mu\alpha\beta\nu} G^{\mu\nu} G^{\alpha\beta} \right) \geq 0, \quad (6)$$

hence (2), (4), (5) and (6) allow to establish the following result:

“If the intrinsic geometry of the spacetime gives:

$$K_2 \neq 0, \quad \left( \frac{R}{24} K_2 + R_{\mu\alpha\beta\nu} G^{\mu\nu} G^{\alpha\beta} \right) \neq 0, \quad (7)$$

then  $R_4$  accepts local and isometric embedding into  $E_5$ ”.

In other words, (7) are intrinsic sufficient conditions for the corresponding embedding.

Let's remember that it is impossible the embedding of any empty  $R_4$  into  $E_5$  [13, 15-17], in fact, if  $R^{\mu\nu} \equiv R^{\tau}{}_{\mu\nu}{}^{\tau}$  then from (1.a):

$$b_\mu{}^a b^{av} = b b_{\mu\nu}, \quad b^{\mu a} b_{\mu a} = b^2, \quad b = b^\nu{}_\nu, \quad (8)$$

therefore:

$$R^{\mu\nu\alpha\beta} R_{\mu\tau\alpha\gamma} = \varepsilon b^2 R^\nu_{\gamma\tau}{}^\beta{}_\alpha \quad (9)$$

which allows to calculate the Synge's gravitational density [18, 19]:

$$F \equiv \frac{1}{2} (R_{\mu\nu} t^\mu t^\nu)^2 + R^{\mu\nu\alpha\beta} R_{\mu\tau\alpha\gamma} t_\nu t_\beta t^\tau t^\gamma = 0 \quad (10)$$

where  $t^\mu$  is an arbitrary timelike vector, hence  $R_4$  is the Minkowski spacetime; thus any vacuum metric has an embedding class  $\geq 2$ .

**Thompson-schrank's Theorem:** Thompson-Schrank [20] obtained the interesting result:

$$R_{\mu\nu} R^{\alpha\nu} = 0 \quad \Rightarrow \quad R \equiv R^\nu_\nu = 0, \quad (11)$$

where  $R_{\alpha\beta}$  and  $R$  are the Ricci tensor and the scalar curvature, respectively. In fact, (11) is equivalent to:

$$\begin{aligned} E_{\mu\nu} E^{\alpha\nu} + \frac{1}{2} R E_\mu{}^\alpha + \frac{1}{4} R^2 \delta_\mu^\alpha &= 0, \\ E_{\alpha\beta} &= R_{\alpha\beta} - \frac{1}{4} R g_{\alpha\beta} \end{aligned} \quad (12)$$

implies  $R = 0$ . If we compare (12) with each Churchill-Plebański type [21-25], we find compatibility only with the types  $[4T]_{[1]}$  and  $[4N]_{[2]}$  for  $R = 0$  thus (11) is immediate. Here we apply (11) to spacetimes of class one to deduce a constraint for the corresponding second fundamental form.

In fact, from (1.a):

$$R_\mu{}^\alpha = \varepsilon (b_{\mu\nu} b^{\alpha\nu} - b b_\mu{}^\alpha), \quad (13)$$

then we can see that the property:

$$b_{\mu\nu} b^{\alpha\nu} = 0, \quad (14)$$

is impossible for 4-spaces of class one because (14) and the Thompson-Schrank's theorem [20] imply  $b = 0$ , thus (13) gives  $R_{\mu\alpha} = 0$ , but we know that no empty spacetime accepts embedding into  $E_5$ . Therefore:

$$b_{\alpha\beta} b^{\beta\gamma} \neq 0, \quad (15)$$

is a constraint for all spacetimes of class one.

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