

## Numerical Solutions of the Linear Volterra Integro-differential Equations: Homotopy Perturbation Method and Finite Difference Method

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**Abstract:** In the research, special type of linear volterra integro-differential equations is considered. This paper compares the Homotopy perturbation method (HPM) with finite difference method for solving these equations. HPM is an analytical procedure for finding the solutions of problems which is based on the constructing a Homotopy with an imbedding parameter  $p$  that is considered as a small parameter. The finite difference method, based upon Simpson rule and Trapezoidal rule, transforms the volterra integro-differential equation into a matrix equation. The results of applying these methods to the linear integro-differential equation show the simplicity and efficiency of these methods.

**Key words:** Volterra integro-differential equations . homotopy perturbation method . finite difference method

### INTRODUCTION

Mathematical modeling of real-life problems usually results in functional equations, e.g. partial differential equations, integral and integro-differential equation, stochastic equations and others. Many mathematical formulations of physical phenomena contain integro-differential equations, these equations arise in fluid dynamics, biological models and chemical kinetics. Integro-differential equations are usually difficult to solve analytically so it is required to obtain an efficient approximate solution. Several numerical methods for approximating the Fredholm or volterra integro-differential equations are known. Single-term Walsh series method for volterra integro-differential equations has been proposed by Sepehrian and Razzaghi [1]. In [2], Brunner applied a collocation-type method to Volterra-Hammerstein integral equation as well as integro-differential equations. Compact finite difference method has been used for integro-differential equations by Zhao and Corless [3]. For methods using a quadrature rule, degenerate kernels, interpolation or extrapolation [4-7]. In Refs [8-10], Taylor series, Chebyshev collocation and Wavelet-Galerkin methods are used for solving such problems. In recent years, the application of homotopy perturbation method (HPM) [11-13] in nonlinear problems has been developed by scientists and engineers, because this method deforms the difficult problem under study into a simple problem which is easy to solve. Most perturbation methods

assume a small parameter exists, but most nonlinear problems have no small parameter at all. Many new methods, such as the variational method [14-16], variational iterations method [17-22], various modified Lindstedt-Poincare methods [23-26] and others [27, 28] are proposed to eliminate the shortcoming arising in the small parameter assumption. A review of recently developed nonlinear analysis methods can be found in [29]. Recently, the applications of homotopy perturbation theory have appeared in the works of many scientist [30-35]; it has become a powerful mathematical tool [36, 37]. In this paper, we propose the use of HPM to solve special type of linear volterra integro-differential equations of the form:

$$\begin{cases} y'(x) + \mu(x)y(x) = f(x) + \lambda \int_a^x k(x,t)y(t)dt, a \leq x \leq b \\ y(a) = y_0 \end{cases} \quad (1)$$

And comparisons are made between finite difference method and homotopy perturbation method. Where the functions  $f(x)$ ,  $\mu(x)$  and the kernel  $k(x, t)$  are known and  $y(x)$  is the solution to be determined.

### HOMOTOPY PERTURBATION METHOD

To illustrate the basic ideas of the homotopy perturbation method, we consider the following nonlinear differential equation:

$$A(y) - f(r) = 0, \quad r \in \Omega, \quad (2)$$

with the boundary conditions

$$B\left(y, \frac{\partial y}{\partial n}\right) = 0, \quad r \in \Omega \quad (3)$$

where  $A$  is a general differential operator,  $B$  is a boundary operator,  $f(r)$  is a known analytical function and  $\Gamma$  is the boundary of the domain  $\Omega$ . Generally speaking, the operator  $A$  can be divided into two parts which are  $L$  and  $N$  where  $L$  is linear, but  $N$  is nonlinear. Therefore equation (2) can therefore be rewritten as follows:

$$L(y) + N(y) - f(r) = 0 \quad (4)$$

By the homotopy perturbation technique, we construct a homotopy  $v(r, p): \Omega \times [0, 1] \rightarrow \mathbb{R}$  which satisfies:

$$H(v, p) = (1-p)[L(v) - L(y_0)] + p[A(v) - f(r)] = 0, \\ p \in [0, 1], \quad r \in \Omega$$

where  $p \in [0, 1]$  is an embedding parameter and  $y_0$  is an initial approximation of equation (2). Obviously, from these definitions we will have:

$$H(v, 0) = L(v) - L(y_0) = 0$$

$$H(v, 1) = A(v) - f(r) = 0$$

The changing process of  $p$  from zero to unity is just that of  $v(r, p)$  from  $y_0(r)$  to  $y(r)$ . In topology, this is called deformation and  $L(v) - L(y_0)$  and  $A(v) - f(r)$  are called homotopy. According to the HPM, we can first use the embedding parameter  $p$  as a "small parameter" and assume that the solution of (5) can be written as a power series in  $p$ :

$$v = v_0 + pv_1 + p^2v_2 + \dots$$

Setting  $p = 1$ , results in the approximate solution of (2):

$$y = \lim_{p \rightarrow 1} v = v_0 + v_1 + v_2 + \dots$$

$$\int_a^{x_1} k(x, t) y(t) dt \approx h \left[ \frac{1}{2} k(x_1, t_0) y(t_0) + \frac{1}{2} k(x_1, t_1) y(t_1) \right] \quad (\text{T. rule})$$

$$\int_a^{x_2} k(x, t) y(t) dt \approx \frac{h}{3} \left[ k(x_2, t_0) y(t_0) + 4 k(x_2, t_1) y(t_1) + k(x_2, t_2) y(t_2) \right] \quad (\text{S. rule})$$

$$\int_a^{x_3} k(x, t) y(t) dt \approx h \left[ \frac{1}{2} k(x_3, t_0) y(t_0) + k(x_3, t_1) y(t_1) + k(x_3, t_2) y(t_2) + \frac{1}{2} k(x_3, t_3) y(t_3) \right] \quad (\text{T. rule})$$

$$\int_a^{x_4} k(x, t) y(t) dt \approx \frac{h}{3} \left[ k(x_4, t_0) y(t_0) + 4 k(x_4, t_1) y(t_1) + 2 k(x_4, t_2) y(t_2) + 4 k(x_4, t_3) y(t_3) + k(x_4, t_4) y(t_4) \right] \quad (\text{S. rule})$$

In order to solve the equation (1) using HPM, we construct the following homotopy:

$$H(v, p) = v' - y'_0 + p y'_0 \\ - p \left[ -\mu(x) v(x) + \lambda \int_a^x k(x, t) v(t) dt + f(x) \right] = 0$$

Substituting (6) in (7) and equating the coefficients of like powers of  $p$ , yield

$$p^0: v'_0 - y'_0 = 0 \quad (8)$$

$$p^1: v'_1 + y'_0 + \mu(x) v_0(x) - \lambda \int_a^x k(x, t) v_0(t) dt - f(x) = 0 \quad (9)$$

$$p^n: v'_n + \mu(x) v_{n-1}(x) - \lambda \int_a^x k(x, t) v_{n-1}(t) dt = 0, \quad n \geq 2 \quad (10)$$

Then starting with an initial approximation  $y_0$  and solving the above equations, we can identify  $v_n$  for  $n = 1, 2, \dots$  and therefore we obtain the  $n$ -th approximation of the exact solution as  $y_n = v_0 + v_1 + \dots + v_n$ .

Note: In this section, we consider  $y(a) = y_0 = \beta$

## FINITE DIFFERENCE METHOD

In this section, we consider volterra integro-differential equation in (1) and approximate to solution by numerical integration and numerical differentiation.

We will subdivide the interval of integration  $(a, x)$  into  $N = 2M$  equal subinterval of with

$$h = \frac{x_N - a}{N}, \quad N \geq 1$$

where  $x_N$  is the end point we choose for  $x$ . We shall set  $t_0 = a$  and  $t_j = t_0 + jh$ . Since we will be using either  $t$  or  $x$  as the independent variable for the solution  $y$ . We will call  $x_0 = t_0 = a$ ,  $x = x_N = t_N$  and  $x_i = a + ih = t_i$ . We will refer to the value of the functions  $f(x)$  and  $p(x)$  at  $x_i$  as  $f(x_i) = f_i$  and  $\mu(x_i) = \mu_i$ , the value of kernel  $k(x, t)$  at  $(x_i, t_j)$  as  $k(x_i, t_j) = k_{ij}$  and the approximate value of the solution  $y(x)$  at  $x_i$  or  $t_i$  as  $y(t_i) = y(x_i) = y_i$  and  $y'(x_i) = y'_i$ .  $k(x_i, t_j)$  Clearly vanishes for  $t_j > x_i$  as the integration ends at  $t_j \leq x_i$ . Note that the particular value  $y(x_0) = y_0$  according to (1). So if we use the trapezoidal rule and Simpson rule with  $n$  subinterval to approximate the integral in the volterra integro-differential equation (1), we have

(11)

$$\begin{aligned} & \vdots \\ & \vdots \\ & \vdots \\ & \int_a^{x_N-1} k(x, t) y(t) dt \approx h \left[ \frac{1}{2} k(x_{N-1}, t_0) y(t_0) + k(x_{N-1}, t_1) y(t_1) + \dots + k(x_{N-1}, t_{N-2}) y(t_{N-2}) + \frac{1}{2} k(x_{N-1}, t_{N-1}) y(t_{N-1}) \right] \quad (\text{T. rule}) \\ & \int_a^{x_N} k(x, t) y(t) dt \approx \frac{h}{3} \left[ k(x_N, t_0) y(t_0) + 4 k(x_N, t_1) y(t_1) + 2 k(x_N, t_2) y(t_2) + \dots + 4 k(x_N, t_{N-1}) y(t_{N-1}) + k(x_N, t_N) y(t_N) \right] \quad (\text{S. rule}) \end{aligned}$$

Also the integro-differential equation (1) is approximated by (compact form)

$$\begin{aligned} y'_1 + \mu_1 y_1 &= f_1 + \lambda h \left[ \frac{1}{2} k_{10} y_0 + \frac{1}{2} k_{11} y_1 \right] \\ y'_2 + \mu_2 y_2 &= f_2 + \frac{\lambda h}{3} [k_{20} y_0 + 4 k_{21} y_1 + k_{22} y_2] \\ y'_3 + \mu_3 y_3 &= f_3 + \lambda h \left[ \frac{1}{2} k_{30} y_0 + k_{31} y_1 + k_{32} y_2 + \frac{1}{2} k_{33} y_3 \right] \\ y'_4 + \mu_4 y_4 &= f_4 + \frac{\lambda h}{3} [k_{40} y_0 + 4 k_{41} y_1 + 2 k_{42} y_2 + 4 k_{43} y_3 + k_{44} y_4] \\ & \vdots \\ & \vdots \\ & \vdots \\ y'_{N-1} + \mu_{N-1} y_{N-1} &= f_{N-1} + \lambda h \left[ \frac{1}{2} k_{N-1,0} y_0 + k_{N-1,1} y_1 + \dots + k_{N-1,N-2} y_{N-2} + \frac{1}{2} k_{N-1,N-1} y_{N-1} \right] \\ y'_N + \mu_N y_N &= f_N + \frac{\lambda h}{3} [k_{N0} y_0 + 4 k_{N1} y_1 + 2 k_{N2} y_2 + \dots + 4 k_{N,N-1} y_{N-1} + k_{NN} y_N] \end{aligned} \quad (12)$$

Now we take advantage of finite differentiation to get

$$\begin{aligned} \frac{y_2 - y_0}{2h} + \mu_1 y_1 &= f_1 + \lambda h \left[ \frac{1}{2} k_{10} y_0 + \frac{1}{2} k_{11} y_1 \right] \\ \frac{y_3 - y_1}{2h} + \mu_2 y_2 &= f_2 + \frac{\lambda h}{3} [k_{20} y_0 + 4 k_{21} y_1 + k_{22} y_2] \\ \frac{y_4 - y_2}{2h} + \mu_3 y_3 &= f_3 + \lambda h \left[ \frac{1}{2} k_{30} y_0 + k_{31} y_1 + k_{32} y_2 + \frac{1}{2} k_{33} y_3 \right] \\ \frac{y_5 - y_3}{2h} + \mu_4 y_4 &= f_4 + \frac{\lambda h}{3} [k_{40} y_0 + 4 k_{41} y_1 + 2 k_{42} y_2 + 4 k_{43} y_3 + k_{44} y_4] \\ & \vdots \\ & \vdots \\ & \vdots \\ \frac{y_N - y_{N-2}}{2h} + \mu_{N-1} y_{N-1} &= f_{N-1} + \lambda h \left[ \frac{1}{2} k_{N-1,0} y_0 + k_{N-1,1} y_1 + \dots + k_{N-1,N-2} y_{N-2} + \frac{1}{2} k_{N-1,N-1} y_{N-1} \right] \\ \frac{3y_N - 4y_{N-1} + y_{N-2}}{2h} + \mu_N y_N &= f_N + \frac{\lambda h}{3} [k_{N0} y_0 + 4 k_{N1} y_1 + 2 k_{N2} y_2 + \dots + 4 k_{N,N-1} y_{N-1} + k_{NN} y_N] \end{aligned} \quad (13)$$

The system (13) consists of N equations and can be written in the following matrix form  $KY = F$ , where

$$K = \begin{pmatrix} -\left(\lambda h^2 k_{11} - 2h\mu_1\right) & 1 & 0 & 0 & \dots & 0 \\ -\left(\frac{8\lambda h^2}{3} k_{21} + 1\right) & -\left(\frac{2\lambda h^2}{3} k_{22} - 2h\mu_2\right) & 1 & 0 & 0 & \dots & 0 \\ -2\lambda h^2 k_{31} - (2\lambda h^2 k_{32} + 1) & -(\lambda h^2 k_{33} - 2h\mu_3) & 1 & 0 & 0 & \dots & 0 \\ -\frac{8\lambda h^2}{3} k_{41} - \frac{4\lambda h^2}{3} k_{42} - \left(\frac{8\lambda h^2}{3} k_{43} + 1\right) & -\left(\frac{2\lambda h^2}{3} k_{44} - 2h\mu_4\right) & 1 & 0 & 0 & \dots & 0 \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ -2\lambda h^2 k_{N-1,1} \dots -2\lambda h^2 k_{N-1,N-3} & -\left(2\lambda h^2 k_{N-1,N-2} + 1\right) & -\left(\lambda h^2 k_{N-1,N-1} - 2h\mu_{N-1}\right) & 1 & \dots & \dots & 0 \\ -\frac{8\lambda h^2}{3} k_{N1} \dots -\frac{8\lambda h^2}{3} k_{N,N-3} & -\left(\frac{4\lambda h^2}{3} k_{N,N-2} + 1\right) & -\left(\frac{8\lambda h^2}{3} k_{N,N-1} + 4\right) & -\left(\frac{2\lambda h^2}{3} k_{NN} - 3 - 2h\mu_N\right) & \dots & \dots & 0 \end{pmatrix}$$

$$Y = \begin{pmatrix} y_1 \\ y_2 \\ \vdots \\ y_N \end{pmatrix}$$

and

$$F = \begin{pmatrix} 2hf_1 + (\lambda h^2 k_{10} + 1)y_0 \\ 2hf_2 + \frac{2\lambda h^2}{3}k_{20}y_0 \\ 2hf_3 + \lambda h^2 k_{30}y_0 \\ 2hf_2 + \frac{2\lambda h^2}{3}k_{40}y_0 \\ \vdots \\ 2hf_{N-1} + \lambda h^2 k_{N-1,0}y_0 \\ 2hf_N + \frac{2\lambda h^2}{3}k_{N0}y_0 \end{pmatrix}$$

### ILLUSTRATIVE EXAMPLES

Now we apply the methods presented to solve the following examples:

**Example 1:** Consider the Volterra integro-differential equation:

$$\begin{cases} y'(x) + y(x) = (x^2 + 2x + 1)e^{-x} + 5x^2 + 8 - \int_0^x ty(t)dt, 0 \leq x \leq 1 \\ y(0) = 10 \end{cases} \quad (14)$$

which has the exact solution  $y(x) = 10 - xe^{-x}$ . The numerical results are represented in Table 1.

To apply the homotopy perturbation method to this equation, we consider  $y_0(x) = 10$  as initial approximation of the exact solution and regarding (8) we start with  $v_0(x) = y_0(x)$ . Since  $v_0(a) = \beta$  and  $y = v_0 + v_1 + v_2 + \dots$  we can set  $v_n(a) = 0$ , ( $n \geq 1$ ) as initial conditions for equations (9) and (10).

**Example 2:** As the second example consider the Volterra integro-differential equation: [1]

$$\begin{cases} y'(x) + y(x) = 1 + 2x + \int_0^x x(1+2x)e^{t(x-t)}y(t)dt, 0 \leq x \leq 1 \\ y(0) = 1 \end{cases} \quad (15)$$

Table 1: Numerical results for example 1

$x_i$	Finite difference method with N = 14	HPM with N = 4	Exact solutions
0.0714	9.933780913	9.933495	9.933495516
0.1428	9.876458584	9.876160	9.876160300
0.2142	9.827589590	9.827049	9.827046197
0.2857	9.785803615	9.785307	9.785292202
0.3571	9.750832589	9.750163	9.750116951
0.4286	9.721466032	9.720923	9.720811832
0.5000	9.697552931	9.696975	9.696734670
0.5714	9.678042251	9.677777	9.677303930
0.6428	9.662857435	9.662846	9.661993413
0.7143	9.651100634	9.651781	9.650327386
0.7857	9.642739377	9.644241	9.641876128
0.8571	9.637018786	9.639963	9.636251847
0.9286	9.633929509	9.638737	9.633104936
1.0000	9.632846912	9.640444	9.632120559

Table 2: Numerical results for example 2

$x_i$	Finite difference method with N = 12	HPM with N = 2	Exact solutions
0.0833	1.00616079	1.0069712	1.006968613
0.1667	1.02810545	1.0282220	1.028167177
0.2500	1.06328807	1.0648317	1.064494459
0.3333	1.11700723	1.1187838	1.117519069
0.4167	1.18751936	1.1931987	1.189592856
0.5000	1.28254044	1.2926690	1.284025417
0.5833	1.40167958	1.4237286	1.405337908
0.6667	1.55629211	1.5955060	1.559623498
0.7500	1.74852146	1.8206228	1.755054657
0.8333	1.99578448	2.1164213	2.002596211
0.9167	2.30508217	2.5066179	2.317010501
1.0000	2.70473683	3.0235181	2.718281828

Table 3: Numerical results for Example 3

$x_i$	Finite difference method with N = 12	HPM with N = 10	Exact solutions
0.0833	0.9409611903	0.9232408643	0.9232408623
0.1667	0.8560967849	0.8582656556	0.8582656554
0.2500	0.8221926661	0.8032653310	0.8032653300
0.3333	0.7521848166	0.7567085596	0.7567085597
0.4167	0.7379172201	0.7172991016	0.7172991041
0.5000	0.6768034806	0.6839397209	0.6839397204
0.5833	0.6667601409	0.6557016132	0.6557016122
0.6667	0.6235899136	0.6317985685	0.6317985690
0.7500	0.6149784301	0.6115650814	0.6115650802
0.8333	0.5862745171	0.5944378021	0.5944378015
0.9167	0.5770459174	0.5799398782	0.5799398730
1.0000	0.5643959592	0.5676676511	0.5676676417

With the exact solution  $y(x) = e^{x^2}$ . Table 2 illustrate the numerical results.

If we want to solve this equation by mean of homotopy perturbation method, Considering  $y_0(x) = 1$  and regarding (8), we start with  $v_0(x) = y_0(x)$  Since  $v_0(a) = \beta$  and  $y = v_0 + v_1 + v_2 + \dots$  we can set  $v_n(a) = 0, (n \geq 1)$  as initial conditions for equations (9) and (10).

**Example 3:** Consider [38]

$$\begin{cases} y'(x) + y(x) = \int_0^x e^{(t-x)} y(t) dt, & 0 \leq x \leq 1 \\ y(0) = 1 \end{cases} \quad (16)$$

with the exact solution  $y(x) = e^{-x} \cosh x$  Results are shown in Table 3.

In order to solve this equation by mean of homotopy perturbation method, We assume  $y_0(x) = 1$  and set  $v_0(x) = y_0(x)$ . We solve the above equations with  $v_0(a) = \beta$  and  $v_n(a) = 0, (n \geq 1)$  as initial conditions.

## CONCLUSION

Integro-differential equations are usually difficult to solve analytically. In many cases, it is required to obtain the approximate solutions. In this work, we proposed the homotopy perturbation method for solving linear volterra integro-differential equations and comparisons were made with the finite difference method. Illustrative examples are included to demonstrate the validity and applicability of these techniques.

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