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On an Inverse Spectral Problem for an Integro-differential Operator

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Abstract: In this paper, we considered an inverse problem with two given spectrum for a boundary value problem with aftereffect and eigenvalue in the boundary condition and we showed that transformation operator was generalized degeneracy and we obtained a new proof of the Hochstadt's theorem concerning the structure of the difference $\tilde{q}(x) - q(x)$.

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INTRODUCTION

We consider boundary value problem with aftereffect on a finite interval and with eigenvalue in the boundary condition:

$$-y'(x) + q(x)y(x) + \int_0^x M(x-t)y(t)dt$$
$$= \lambda^2 y(x), \ 0 \le x \le \pi$$

with the boundary condition

$$y(0,\lambda) - \lambda y(0,\lambda) = 0$$
$$y'(\pi,\lambda) + Hy(\pi,\lambda) = 0$$

Here λ is spectral parameter, $q(x) \in L_2(0,\pi)$ and H are real parameter. The presence of an aftereffect in a mathematical model produces qualitative changes in the study of the inverse problem. In recent years, inverse spectral problems for integro differential operators have been studied by many authors including [1-4].

The generalized degeneracy of transformation operator for Sturm-Liouville and singular Sturm-Liouville operator was showed in [5-7].

In this paper, using Levitan's method, it was shown that kernel K(x,t) of the integral equation is generalized degeneracy.

Now let us consider two boundary value problem with aftereffect on a finite interval and with eigenvalue in the boundary condition:

$$-y'(x) + q(x)y(x) + \int_0^x M(x-t)y(t)dt$$

= $\lambda^2 y(x), 0 \le x \le \pi$ (1.1)

$$y'(0) - \lambda y(0) = 0$$
 (1.2)

$$\mathbf{y}'(\boldsymbol{\pi},\boldsymbol{\lambda}) + \mathbf{H}\mathbf{y}(\boldsymbol{\pi},\boldsymbol{\lambda}) = 0 \tag{1.3}$$

$$-\dot{\mathbf{y}}(\mathbf{x}) + \tilde{\mathbf{q}}(\mathbf{x})\mathbf{y}(\mathbf{x}) + \int_{0}^{x} \widetilde{\mathbf{M}}(\mathbf{x} - \mathbf{t})\mathbf{y}(\mathbf{t})d\mathbf{t}$$

= $\lambda^{2}\mathbf{y}(\mathbf{x}), \ 0 \le \mathbf{x} \le \pi$ (1.4)

$$y(0) - \lambda y(0) = 0$$
$$y(\pi, \lambda) + \widetilde{H}y(\pi, \lambda) = 0$$
(1.5)

where $\tilde{q}(x) \in L_2(0,\pi)$ and \tilde{H} is real parameter.

Denote the spectrum of the first problem by $\{\lambda_n\}_{n\geq 1}$ and the spectrum of the second by $\{\tilde{\lambda}_n\}_{n\geq 1}$. Next, we denote by $\varphi(x,\lambda)$ the solution of (1.1) and by $\tilde{\varphi}(x,\lambda)$ the solution of (1.4), satisfying the initial condition (1.2). The representation

$$\tilde{\phi}(\mathbf{x},\lambda) = \phi(\mathbf{x},\lambda) + \int_0^x K(\mathbf{x},t)\phi(t,\lambda)dt$$
(1.6)

holds, where K(x,t) satisfies the equation

$$K_{tt}(x,t) - K_{xx}(x,t) + (q(x) - \tilde{q}(t))K(x,t)$$

+M(x-t) - $\widetilde{M}(x-t)$ (1.7)
+ $\int_{t}^{x} (M(x-\xi) - \widetilde{M}(x-\xi))K(\xi,t)d\xi = 0$

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and the conditions

$$2\frac{\mathrm{dK}(\mathbf{x},\mathbf{x})}{\mathrm{dx}} = \tilde{q}(\mathbf{x}) - q(\mathbf{x}) \tag{1.8}$$

$$K(x,0)=0$$
 (1.9)

We put

$$\mathbf{c}_{n} = \int_{0}^{\pi} \boldsymbol{\varphi}^{2}(\mathbf{x}, \boldsymbol{\lambda}_{n}) d\mathbf{x}, \quad \tilde{\mathbf{c}}_{n} = \int_{0}^{\pi} \boldsymbol{\varphi}^{2}(\mathbf{x}, \tilde{\boldsymbol{\lambda}}_{n}) d\mathbf{x}$$
(1.10)

$$\rho(\lambda) =_{\lambda_n < \lambda} \frac{1}{c_n}, \quad \tilde{\rho}(\lambda) = \frac{1}{\tilde{\lambda}_n < \lambda} \frac{1}{\tilde{c}_n}$$
(1.11)

The function $\rho(\lambda)$ ($\tilde{\rho}(\lambda)$) is called the spectral function of problem (1.1)-(1.3) [(1,4), (1.5)]. Problem (1.1)-(1.3) is regarded as an unperturbed problem, while (1.4), (1.5) is considered as a perturbation of (1.1)-(1.3).

It is known fact that in [5] the knowledge of two spectra for given Sturm-Liouville equation makes it possible to recover its spectral function, i.e., to find the numbers $\{c_n\}$. More exactly, suppose that, in addition to the spectrum of problem (1.1)-(1.3), we also know the spectrum $\{\mu_n\}$ of the problem

$$-y'(x) + q(x)y(x) + \int_{0}^{x} M(x-t)y(t)dt$$

= $\lambda^{2}y(x), 0 \le x \le \pi$
 $y'(0) - \lambda y(0) = 0$ (1.12)
 $y'(\pi, \lambda) + H_{1}y(\pi, \lambda) = 0, H_{1} \ne H$

Knowing $\{\lambda_n\}$ and $\{\mu_n\}$, we can calculate the numbers $\{c_n\}$. Similarly, for (1.4), if, in addition to $\{\tilde{\lambda}_n\}$, we also know the spectrum $\{\tilde{\mu}_n\}$ determined by the boundary conditions

$$y'(0) - \lambda y(0) = 0$$
$$y'(\pi, \lambda) + \widetilde{H}_1 y(\pi, \lambda) = 0, \widetilde{H}_1 \neq \widetilde{H}$$
(1.13)

then it follows that we can determine the numbers $\{\tilde{c}_n\}$. It is also shown that in [4]

$$\lambda_{n} = n + \frac{1}{\pi} + \frac{H}{n\pi} + \frac{1}{2n\pi} \int_{0}^{\pi} q(t) dt + \frac{k_{n}}{n}, \{k_{n}\} \in I_{2} \text{ and } n \in \mathbb{N} - \{0\}$$
(1.14)

$$\varphi(x,\lambda) = \cos \lambda x + \sin \lambda x + \frac{1}{\lambda} \int_0^x \sin \lambda (x-\tau) \left[q(\tau) \varphi(\tau,\lambda) + \int_0^\tau M(\tau-s) \varphi(s,\lambda) ds \right] d\tau$$
(1.15)

Theorem 1. Consider the operator

$$Ly(x) = -y'(x) + q(x)y(x) + \int_0^x M(x-t)y(t) dt = \lambda^2 y(x), \ 0 \le x \le \pi$$
(1.16)

subject to the boundary conditions

$$y(0) - \lambda y(0) = 0$$
 (1.17)

$$y'(\pi,\lambda) + Hy(\pi,\lambda) = 0$$
(1.18)

Let $\{\lambda_n\}$ be the spectrum of L subject to (1.17) and (1.18). If (1.18) is replaced by the new boundary condition

$$\mathbf{y}'(\boldsymbol{\pi},\boldsymbol{\lambda}) + \mathbf{H}_{1}\mathbf{y}(\boldsymbol{\pi},\boldsymbol{\lambda}) = 0 \tag{1.19}$$

then a new operator and a new spectrum, say $\{\mu_n\}$, result. Now, consider the second operator

$$\widetilde{L}y(x) = -y'(x) + \widetilde{q}(x)y(x) + \int_0^x \widetilde{M}(x-t)y(t)dt = \lambda^2 y(x), \ 0 \le x \le \pi$$
(1.20)

Suppose that \tilde{L} has the spectrum $\{\tilde{\lambda}_n\}$ with $\{\tilde{\lambda}_n\} = \{\lambda_n\}$ for all n under the boundary conditions (1.17) and

$$y'(\pi,\lambda) + \widetilde{H}y(\pi,\lambda) = 0$$
(1.21)
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 \tilde{L} with the boundary conditions (1.17) and

$$y'(\pi,\lambda) + \widetilde{H}_{1}y(\pi,\lambda) = 0$$
(1.22)

is assumed to have the spectrum $\{\widetilde{\mu}_n\}$. Assuming that H, $H_l \neq H$, \widetilde{H} and $\widetilde{H}_l \neq \widetilde{H}$ are real numbers that are not infinite.

Denote by Λ_0 the finite index set for which $\tilde{\mu}_n \neq \mu_n$ and by Λ the infinite index set for which $\tilde{\mu}_n = \mu_n$. Under the above assumptions, it follows that the kernel K(x,t) is degenerate in the extended sense:

$$K(x,t) = \sum_{\Lambda_0} c_n \tilde{\varphi}_n(x) \varphi_n(t)$$
(1.23)

where ϕ_n and $\tilde{\phi}_n$ are suitable solutions of (1.1) and (1.4).

Proof: It follows from (1.6) that

$$\tilde{\varphi}'(x,\lambda) = \varphi'(x,\lambda) + K(x,x)\varphi(x,\lambda) + \int_0^x \frac{\partial K}{\partial x}\varphi(t,\lambda)dt$$
(1.24)

and

$$\tilde{\varphi}'(x,\lambda) + \tilde{H} \,\tilde{\varphi}(x,\lambda) = \varphi'(x,\lambda) + \tilde{H}\varphi(x,\lambda) + K(x,x)\varphi(x,\lambda) + \int_0^x \left(\frac{\partial K}{\partial x} + \tilde{H}K\right)\varphi(t,\lambda)dt \tag{1.25}$$

Substituing $x=\pi$ and $\lambda = \lambda_n$ into the last equation and using the boundary conditions (1.18), we obtain

$$\left(\widetilde{H} - H\right)\phi(\pi,\lambda_{n}) + K(\pi,\pi)\phi(\pi,\lambda_{n}) + \int_{0}^{\pi} \left(\frac{\partial K}{\partial x} + \widetilde{H}K\right)_{x=\pi} \phi(t,\lambda_{n}) dt = 0$$
(1.26)

As $n \to \infty$ and $\varphi(\pi, \lambda_n) \to (-1)^n (\cos 1 + \sin 1)$, the integral on the right-hand side tends to zero. Therefore, from (1.26) we get

$$K(\pi, \dot{\eta} = H - \dot{H}$$
(1.27)

$$\int_{0}^{\pi} \left(\frac{\partial K}{\partial x} + \widetilde{H} K \right)_{x=\pi} \phi(t, \lambda_{n}) dt = 0, n = 0, 1, \dots$$
(1.28)

Since the systems of functions $\varphi(t,\lambda_n)$ is complete, it follows from the last equation that

$$\left(\frac{\partial K}{\partial x} + \widetilde{H}K\right)_{x=\pi} = 0, \ 0 \le t \le \pi$$
(1.29)

We now use the equation imposed on the second mentioned spectrum. Using (1.6) again, we obtain

$$\tilde{\varphi}'(x,\lambda) + \tilde{H}_{1}\tilde{\varphi}(x,\lambda) = \varphi'(x,\lambda) + \tilde{H}_{1}\varphi(x,\lambda) + K(x,x)\varphi(x,\lambda) + \int_{0}^{x} \left(\frac{\partial K}{\partial x} + \tilde{H}_{1}K\right)\varphi(t,\lambda)dt$$
(1.30)

Setting $x=\pi$ and $\lambda = \mu_n$ ($n \in \Lambda$) and using (1.19), we get

$$\int_{0}^{\pi} \left(\frac{\partial K}{\partial x} + \widetilde{H}_{1} K \right)_{x=\pi} \phi(t, \mu_{n}) dt + \left(\widetilde{H}_{1} - H_{1} \right) \phi(\pi, \mu_{n}) + K(\pi, \pi) \phi(\pi, \mu_{n}) = 0$$
(1.31)

In the last equation as $n \rightarrow \infty$, the left-hand side tends to zero and $\varphi(\pi, \mu_n) \rightarrow (-1)^n (\cos 1 + \sin 1)$. Therefore;

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$$K(\pi, \eta) = H_1 - \widetilde{H}_1$$
(1.32)

$$\int_{0}^{\pi} \left(\frac{\partial K}{\partial x} + \widetilde{H}_{1} K \right)_{x=\pi} \phi(t, \mu_{n}) dt = 0, n \in \Lambda$$
(1.33)

Comparing (1.27) and (1.32), we obtain $H - \tilde{H} = H_1 - \tilde{H}_1$. For $n \in \Lambda_0$, relation (1.30) for $(x = \pi \text{ and } \lambda = \mu_n)$ yields

$$\int_{0}^{\pi} \left(\frac{\partial K}{\partial x} + \widetilde{H}_{1} K \right)_{x=\pi} \phi(t,\mu_{n}) dt = \widetilde{\phi}'(\pi,\mu_{n}) + \widetilde{H}_{1} \widetilde{\phi}(\pi,\mu_{n})$$
(1.34)

It follows from (1.33) and (1.34) that

$$\left(\frac{\partial K}{\partial x} + \widetilde{H}_{1}K\right)_{x=\pi} = \sum_{\Lambda_{0}} \frac{\widetilde{\phi}'(\pi, \mu_{n}) + \widetilde{H}_{1}\widetilde{\phi}(\pi, \mu_{n})}{\left\|\phi(t, \mu_{n})\right\|^{2}} \phi(t, \mu_{n}), 0 \le t \le \pi$$
(1.35)

We derive from (1.29) and (1.35) the following equations:

$$K(\pi,t) = \frac{1}{\widetilde{H}_{1} - \widetilde{H}} \sum_{\Lambda_{0}} \frac{\widetilde{\phi}(\pi,\mu_{n}) + \widetilde{H}_{1} \widetilde{\phi}(\pi,\mu_{n})}{\left\|\phi(t,\mu_{n})\right\|^{2}} \phi(t,\mu_{n})$$
(1.36)

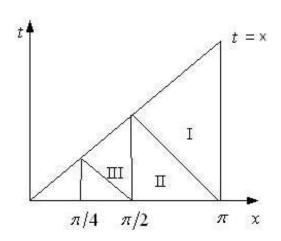
$$\left(\frac{\partial K(x,t)}{\partial x}\right)_{x=\pi} = -\frac{\widetilde{H}}{\widetilde{H}_{1} - \widetilde{H}} \sum_{\Lambda_{0}} \frac{\widetilde{\phi}'(\pi,\mu_{n}) + \widetilde{H}_{1}\widetilde{\phi}(\pi,\mu_{n})}{\left\|\phi(t,\mu_{n})\right\|^{2}} \phi(t,\mu_{n}), 0 \le t \le \pi$$
(1.37)

The function K(x,t) satisfies (1.7). Therefore, it follows from the initial conditions (1.36) and (1.37) that, in triangle I (Fig. 1), we have

$$K(x,t) = \frac{1}{\widetilde{H}_{1} - \widetilde{H}} \sum_{\Lambda_{0}} \frac{\widetilde{\phi}(\pi,\mu_{n}) + \widetilde{H}_{1} \widetilde{\phi}(\pi,\mu_{n})}{\left\|\phi(t,\mu_{n})\right\|^{2}} \left[\widetilde{c}\left(x,\mu_{n}\right) - \widetilde{H}\widetilde{t}(x,\mu_{n})\right] \phi(t,\mu_{n})$$
(1.38)

where $\tilde{c}(x,\lambda)$ and $\tilde{t}(x,\lambda)$ are the solutions of (1.4) satisfying the initial conditions

$$\tilde{\mathbf{c}}(\pi,\lambda) = \tilde{\mathbf{t}}'(\pi,\lambda) = 1, \quad \tilde{\mathbf{c}}'(\pi,\lambda) = \tilde{\mathbf{t}}(\pi,\lambda) = 0$$
(1.39)



The function K(x,t) and sum (1.38) satisfy (1.9); therefore, they coincide in triangle II; consequently, they coincide in triangle III, because solutions of (1.7) satisfy the same initial conditions on the line $x = \frac{\pi}{2}$, etc., i.e., K(x,t) is expressed by (1.38) throughout the triangle $0 \le x \le t \le \pi$.

Theorem 2: If the spectra and $\{\lambda_n\}$ and $\{\tilde{\lambda}_n\}$ coincide and $\{\mu_n\}$ and $\{\tilde{\mu}_n\}$ differ in a finite number of their terms, i.e., $\tilde{\mu}_n = \mu_n$ for $n \in \Lambda$, then

$$\tilde{q}(x) - q(x) = \sum_{\Lambda_0} \tilde{c}_n \frac{d}{dx} \left(\tilde{\varphi}_n \cdot \varphi_n \right)$$
(1.40)

where ϕ_n and $\tilde{\phi}_n$ are suitable solutions of (1.1) and (1.4)

Proof: We obtain from (1.8) the equation

$$\tilde{q}(x) - q(x) = 2\frac{dK(x,x)}{dx}$$
(1.41)

Differentiating (1.38) and setting t = x, we obtain

$$\tilde{q}(x) - q(x) = \frac{2}{\tilde{H}_1 - \tilde{H}} \sum_{\Lambda_0} \frac{\tilde{\phi}(\pi, \mu_n) + \tilde{H}_1 \tilde{\phi}(\pi, \mu_n)}{\left\| \phi(t, \mu_n) \right\|^2} \frac{d}{dx} \left\{ \left[\tilde{c}(x, \mu_n) - \tilde{H}\tilde{t}(x, \mu_n) \right] \phi(x, \mu_n) \right\}$$
(1.42)

Consequently,

$$\tilde{q}(x) - q(x) = \sum_{\Lambda_0} \tilde{c}_n \frac{d}{dx} \left(\tilde{\varphi}_n \, \varphi_n \right)$$
(1.43)

where

$$\widetilde{c}(x,\mu_n) - \widetilde{H}\widetilde{t}(x,\mu_n) = \widetilde{\phi}_n, \ \phi(x,\mu_n) = \phi_n(x,\mu_n)$$

and

$$\tilde{\mathbf{c}}_{n} = \frac{2\left[\tilde{\boldsymbol{\phi}}'(\boldsymbol{\pi},\boldsymbol{\mu}_{n}) - \widetilde{\mathbf{H}}_{1}\tilde{\boldsymbol{\phi}}(\boldsymbol{\pi},\boldsymbol{\mu}_{n})\right]}{\left(\widetilde{\mathbf{H}}_{1} - \widetilde{\mathbf{H}}\right)\left\|\boldsymbol{\phi}(\mathbf{t},\boldsymbol{\mu}_{n})\right\|^{2}}$$
(1.44)

This completes the proof of Theorem 2. We note that similar problem was investigated [7].

REFERENCES

- Freiling, G. and V.A. Yurko, 2001. Inverse Sturm-Liouville problems and their applications. NOVA science Publishers, New York.
- Kuryshova Yu, 2007. The inverse spectral problem for integro-differential operators. Mat. Zametki, 81: 855-866; English translation in Math. Notes, 81: 767-777.

- 3. Buterin, S., 2007. On an inverse spectral problem for a convulation integro-differential operator. Result in Mathematics, 50: 173-181.
- Dabbaghian, A.H., S. Akbarpoor and J. Vahidi, 2011. The uniqueness theorem for boundary value problem with aftereffect and eigenvalue in the boundary condition. The journal of Mathematics and Computer Science, 2: 483-487.
- Levitan, B.M., 1978. On the determination of the Sturm-Liouville operator from one and two spectra. Izv. Akad. Nauk SSSR, Ser. Mat., 42: 185-199.
- Panakhov, E.S. and R. Yilmazer, 2006. On inverse problem for singular Sturm-Liouville operator from two spectra. Ukrainian Mathematical Journal, 58: 147-154.
- Hochstadt, H., 1973. The inverse Sturm-Liouville problem. Commun. Pure Appl. Math., 106: 715-729.
- 8. Hochstadt H., 1973. The inverse Sturm-Liouville problem, Commun. Pure Appl. Math. 106 715-729.