

Focal Curve of Timelike Biharmonic Curve in the Lorentzian Heisenberg Group Heis^3

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Abstract: In this paper, we study focal curve of timelike biharmonic curve in the Heisenberg group Heis^3 . We characterize focal curve of timelike biharmonic curve in terms of curvature and torsion of biharmonic curve in the Heisenberg group Heis^3 . Finally, we construct parametric equations of focal curve of timelike biharmonic curve.

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INTRODUCTION

Darboux had found how to determine the evolutes of a curve, that is, the curves whose tangents are normals of γ . Moreover, he had shown that the focal surface of γ is foliated by the evolutes and all of them lie on the focal surface.

The differential geometry of space curves is a classical subject which usually relates geometrical intuition with analysis and topology. For any unit speed curve γ , the focal curve $C\gamma$ is defined as the centers of the osculating spheres of γ . Since the center of any sphere tangent to γ at a point lies on the normal plane to γ at that point, the focal curve of γ may be parameterized using the Frenet frame $(t(s), n_1(s), n_2(s))$ of γ as follows:

$$C_\gamma(s) = (\gamma + c_1 n_1 + c_2 n_2)(s),$$

Where the coefficients c_1, c_2 are smooth functions that are called focal curvatures of γ .

The aim of this paper is to study focal curve of a timelike biharmonic curve in the Lorentzian Heisenberg group Heis^3 .

Harmonic maps $f: (M, g) \rightarrow (N, h)$ between Riemannian manifolds are the critical points of the energy

$$E(f) = \frac{1}{2} \int_M |df|^2 v_g, \quad (1.1)$$

and they are therefore the solutions of the corresponding Euler-Lagrange equation. This equation is given by the vanishing of the tension field

$$\tau(f) = \text{trace} \nabla df \quad (1.2)$$

The bienergy of a map f by

$$E_2(f) = \frac{1}{2} \int_M |\tau(f)|^2 v_g, \quad (1.3)$$

and say that is biharmonic if it is a critical point of the bienergy.

The first and the second variation formula for the bienergy, showing that the Euler-Lagrange equation associated to E_2 is

$$\begin{aligned} \tau_2(f) &= -J^f(\tau(f)) = -\Delta \tau(f) - \text{trace} R^N(df, \tau(f))df \\ &= 0, \end{aligned} \quad (1.4)$$

Where J^f is the Jacobi operator of f . The equation $\tau_2(f) = 0$ is called the biharmonic equation. Since J^f is linear, any harmonic map is biharmonic. Therefore, we are interested in proper biharmonic maps, that is non-harmonic biharmonic maps.

In this paper, we study focal curve of timelike biharmonic curve in the Heisenberg group Heis^3 . We characterize focal curve of timelike biharmonic curve in terms of curvature and torsion of biharmonic curve in the Heisenberg group Heis^3 . Finally, we construct parametric equations of focal curve of timelike biharmonic curve.

The Lorentzian Heisenberg Group Heis^3 : The Lorentzian Heisenberg group Heis^3 can be seen as the space \mathbb{R}^3 endowed with the following multiplication:

$$(\bar{x}, \bar{y}, \bar{z})(x, y, z) = (\bar{x} + x, \bar{y} + y, \bar{z} + z - \bar{x}y + x\bar{y}).$$

$Heis^3$ is a three-dimensional, connected, simply connected and 2-step nilpotent Lie group.

The Lorentz metric g is given by

$$g = dx^2 + dy^2 + (xdy + dz)^2.$$

The Lie algebra of $Heis^3$ has an orthonormal basis

$$\mathbf{e}_1 = \frac{\partial}{\partial z}, \mathbf{e}_2 = \frac{\partial}{\partial y} - x \frac{\partial}{\partial z}, \mathbf{e}_3 = \frac{\partial}{\partial x} \quad (2.1)$$

for which we have the Lie products

$$[\mathbf{e}_2, \mathbf{e}_3] = 2\mathbf{e}_1, [\mathbf{e}_3, \mathbf{e}_1] = 0, [\mathbf{e}_2, \mathbf{e}_1] = 0$$

With

$$g(\mathbf{e}_1, \mathbf{e}_1) = g(\mathbf{e}_2, \mathbf{e}_2) = 1, g(\mathbf{e}_3, \mathbf{e}_3) = -1.$$

Proposition 2.1: For the covariant derivatives of the Levi-Civita connection of the left-invariant metric g , defined above, the following is true:

$$\nabla = \begin{pmatrix} 0 & \mathbf{e}_3 & \mathbf{e}_2 \\ \mathbf{e}_3 & 0 & \mathbf{e}_1 \\ \mathbf{e}_2 & -\mathbf{e}_1 & 0 \end{pmatrix}, \quad (2.2)$$

Where the (i,j) -element in the table above equals $\nabla_{\mathbf{e}_i} \mathbf{e}_j$ for our basis

$$\{\mathbf{e}_k, k = 1, 2, 3\} = \{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}.$$

We adopt the following notation and sign convention for Riemannian curvature operator:

$$R(X, Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]} Z.$$

The Riemannian curvature tensor is given by

$$R(X, Y, Z, W) = -g(R(X, Y)Z, W).$$

Moreover we put

$$R_{abc} = R(\mathbf{e}_a, \mathbf{e}_b)\mathbf{e}_c, R_{abcd} = R(\mathbf{e}_a, \mathbf{e}_b, \mathbf{e}_c, \mathbf{e}_d),$$

Where the indices a, b, c and d take the values 1, 2 and 3.

$$R_{232} = -3R_{131} = -3\mathbf{e}_3,$$

$$R_{133} = -R_{122} = \mathbf{e}_1,$$

$$R_{233} = -3R_{121} = -3\mathbf{e}_2,$$

and

$$R_{1212} = -1, R_{1313} = 1, R_{2323} = -3. \quad (2.3)$$

Timelike Biharmonic Curves in the Lorentzian Heisenberg Group $Heis^3$: Let $\gamma: I \rightarrow Heis^3$ be a timelike curve on the Lorentzian Heisenberg group $Heis^3$ parametrized by arc length. Let $\{t, n_1, n_2\}$ be the Frenet frame fields tangent to the Lorentzian Heisenberg group $Heis^3$ along γ defined as follows:

t is the unit vector field γ' tangent to γ , n_1 is the unit vector field in the direction of $\nabla_t t$ (normal to γ) and n_2 is chosen so that $\{t, n_1, n_2\}$ is a positively oriented orthonormal basis. Then, we have the following Frenet formulas:

$$\begin{aligned} \nabla_t t &= \kappa n_1, \\ \nabla_t n_1 &= \kappa t + \tau n_2, \\ \nabla_t n_2 &= -\tau n_1, \end{aligned} \quad (3.1)$$

Where κ is the curvature of γ and τ is its torsion.

With respect to the orthonormal basis $\{\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3\}$ we can write

$$t = t_1 \mathbf{e}_1 + t_2 \mathbf{e}_2 + t_3 \mathbf{e}_3,$$

$$n_1 = n_1^1 \mathbf{e}_1 + n_1^2 \mathbf{e}_2 + n_1^3 \mathbf{e}_3,$$

$$n_2 = t \times n_1 = n_2^1 \mathbf{e}_1 + n_2^2 \mathbf{e}_2 + n_2^3 \mathbf{e}_3.$$

Theorem 3.1: (see [18]) Let $\gamma: I \rightarrow Heis^3$ be a non-geodesic timelike curve on the Lorentzian Heisenberg group $Heis^3$ parametrized by arc length. γ is a timelike non-geodesic biharmonic curve if and only if

$$\kappa = \text{constant} \neq 0,$$

$$\kappa^2 - \tau^2 = -1 + 4(n_2^1)^2 \quad (3.2)$$

$$\tau' = -2n_1^1 n_2^1.$$

Corollary 3.2: (see [18]) Let $\gamma: I \rightarrow Heis^3$ be a non-geodesic timelike curve on the Lorentzian Heisenberg group $Heis^3$ parametrized by arc length. γ is biharmonic if and only if

$$\begin{aligned}\kappa &= \text{constant} \neq 0, \\ \tau &= \text{constant}, \\ n_1^1 n_2^1 &= 0, \\ \kappa^2 - \tau^2 &= -1 + 4(n_2^1)^2.\end{aligned}\quad (3.3)$$

Theorem 3.3.: (see [18]) Let $\gamma : I \rightarrow Heis^3$ be a non-geodesic timelike curve on Lorentzian Heisenberg group $Heis^3$ parametrized by arc length. If $n_1^1 \neq 0$, then γ is not biharmonic.

Theorem 3.4.: (see [18]) Let $\gamma : I \rightarrow Heis^3$ be a non-geodesic timelike biharmonic curve on the Lorentzian Heisenberg group $Heis^3$ parametrized by arc length. If $N_1 = 0$, then

$$\mathbf{t}(s) = \sinh \phi_0 \mathbf{e}_1 + \cosh \phi_0 \sinh \psi(s) \mathbf{e}_2 + \cosh \phi_0 \cosh \psi(s) \mathbf{e}_3, \quad (3.4)$$

Where $\phi_0 \in \mathbb{R}$.

Focal Curve of Timelike Biharmonic Curve in the Lorentzian Heisenberg Group $Heis^3$: For a unit speed curve γ , the curve consisting of the centers of the osculating spheres of γ is called the parametrized focal curve of γ . The hyperplanes normal to γ at a point consist of the set of centers of all spheres tangent to γ at that point. Hence the center of the osculating spheres at that point lies in such a normal plane. Therefore, denoting the focal curve by C_γ , we can write

$$C_\gamma(s) = (\gamma + c_1 \mathbf{n} + c_2 \mathbf{b})(s), \quad (4.1)$$

Where the coefficients c_1, c_2 are smooth functions of the parameter of the curve γ , called the first and second focal curvatures of γ , respectively. Further, the focal curvatures c_1, c_2 are defined by

$$c_1 = \frac{1}{\kappa}, c_2 = \frac{c_1'}{\tau}, \kappa \neq 0, \tau \neq 0. \quad (4.2)$$

Lemma 4.1.: Let $\gamma : I \rightarrow Heis^3$ be a unit speed timelike biharmonic curve and C_γ its focal curve on $Heis^3$. Then,

$$c_1 = \frac{1}{\kappa} = \text{constant and } c_2 = 0. \quad (4.3)$$

Proof: Using (3.3) and (4.2), we get (4.3).

Lemma 4.2. Let $\gamma : I \rightarrow Heis^3$ be a unit speed timelike biharmonic curve and C_γ its focal curve on $Heis^3$. Then,

$$C_\gamma(s) = (\gamma + c_1 \mathbf{n})(s). \quad (4.4)$$

Theorem 4.3.: Let $\gamma : I \rightarrow Heis^3$ be a unit speed timelike biharmonic curve and C_γ its focal curve on $Heis^3$. Then, the parametric equations of C_γ are

$$\begin{aligned}x_{C_\gamma}(s) &= \frac{c_1}{\kappa} (\Re \cosh \phi_0 \sinh[\Re s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\Re s + \varepsilon]) \\ &\quad + \frac{1}{\Re} \cosh \phi_0 \sinh[\Re s + \varepsilon] + \varepsilon_1, \\ y_{C_\gamma}(s) &= \frac{c_1}{\kappa} (\Re \cosh \phi_0 \cosh[\Re s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\Re s + \varepsilon]) \\ &\quad + \frac{1}{\Re} \cosh \phi_0 \sinh[\Re s + \rho] + \varepsilon_2, \\ z_{C_\gamma}(s) &= \frac{c_1}{\kappa} (\Re \cosh \phi_0 \cosh[\Re s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\Re s + \varepsilon]) \\ &\quad + \left(\frac{\kappa}{\Re} \cosh \phi_0 \sinh[\Re s + \varepsilon] + \frac{2\kappa}{\Re^2} \sinh \phi_0 \cosh \phi_0 \cosh[\Re s + \varepsilon] + \varepsilon_3 s + \varepsilon_4 \right) \\ &\quad + \frac{1}{\Re} \cosh \phi_0 \sinh[\Re s + \varepsilon] - \frac{1}{\Re} \cosh^2 \phi_0 \left(-\frac{s}{2} + \frac{\sinh 2[\Re s + \varepsilon]}{4\Re} \right) \\ &\quad - \frac{\varepsilon_1}{\Re} \cosh \phi_0 \cosh[\Re s + \varepsilon] + \varepsilon_5,\end{aligned}\quad (4.5)$$

Where $\varepsilon, \varepsilon_1, \varepsilon_2, \varepsilon_3, \varepsilon_4, \varepsilon_5$, are constants of integration and $\Re = (\pm \frac{\kappa}{\cosh \phi_0} - 2 \sinh \phi_0)$.

Proof: The covariant derivative of the vector field \mathbf{t} is:

$$\nabla_{\mathbf{t}} \mathbf{t} = t_1' \mathbf{e}_1 + (t_2' + 2t_1 t_3) \mathbf{e}_2 + (t_3' + 2t_1 t_2) \mathbf{e}_3. \quad (4.6)$$

From (3.4), we have

$$\begin{aligned}\nabla_{\mathbf{t}} \mathbf{t} &= (\psi' \cosh \phi_0 \cosh \psi(s) + 2 \sinh \phi_0 \cosh \phi_0 \cosh \psi(s)) \mathbf{e}_2 \\ &\quad + (\psi' \cosh \phi_0 \sinh \psi(s) + 2 \sinh \phi_0 \cosh \phi_0 \cosh \psi(s)) \mathbf{e}_3.\end{aligned}\quad (4.7)$$

Since $|\nabla_{\mathbf{t}} \mathbf{t}| = \kappa$ we obtain

$$\psi(s) = (\pm \frac{\kappa}{\cosh \phi_0} - 2 \sinh \phi_0) s + \varepsilon, \quad (4.8)$$

Where $\varepsilon \in \mathbb{R}$.

Thus (3.4) and (4.8), imply

$$\mathbf{t} = \sinh \phi_0 \mathbf{e}_1 + \cosh \phi_0 \sinh[\Re s + \varepsilon] \mathbf{e}_2 + \cosh \phi_0 \cosh[\Re s + \varepsilon] \mathbf{e}_3, \quad (4.9)$$

Where $\Re = (\pm \frac{\kappa}{\cosh \phi_0} - 2 \sinh \phi_0)$.

Using (3.1) in (4.9), we obtain

$$\mathbf{t} = (\cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon], \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon], \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \\ - x(s) \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon].$$

From (2.1), we get

$$\mathbf{t} = (\cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon], \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon], \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \\ - \left(\frac{1}{\mathfrak{R}} \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon] + \varepsilon_1 \right) \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon], \quad (4.10)$$

Where ε_1 is constant of integration.

From (3.1) and (4.9), we get

$$\nabla_{\mathbf{t}} \mathbf{t} = (\mathfrak{R} \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \mathbf{e}_2 \\ + (\mathfrak{R} \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \mathbf{e}_3.$$

Where $\mathfrak{R} = (\pm \frac{\kappa}{\cosh \phi_0} - 2 \sinh \phi_0)$.

By the use of Frenet formulas, we get

$$\mathbf{n}_1 = \frac{1}{\kappa} \nabla_{\mathbf{t}} \mathbf{t} \\ = \frac{1}{\kappa} [(\mathfrak{R} \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \mathbf{e}_2 \\ + (\mathfrak{R} \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \mathbf{e}_3]. \quad (4.11)$$

Substituting (2.1) in (4.11), we have

$$\mathbf{n}_1 = \frac{1}{\kappa} ((\mathfrak{R} \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]), \\ (\mathfrak{R} \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]), \\ (\mathfrak{R} \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon] + 2 \sinh \phi_0 \cosh \phi_0 \cosh[\mathfrak{R}s + \varepsilon]) \quad (4.12)$$

$$. (\frac{\kappa}{\mathfrak{R}} \cosh \phi_0 \sinh[\mathfrak{R}s + \varepsilon] + \frac{2\kappa}{\mathfrak{R}^2} \sinh \phi_0 \cosh \phi_0 \\ \cosh[\mathfrak{R}s + \varepsilon] + \varepsilon_3 + \varepsilon_4)),$$

Where $\varepsilon_3, \varepsilon_4$ are constants of integration.

Next, we substitute (4.10) and (4.12) into (4.4), we get (4.5). The proof is completed.

Using Mathematica in Theorem 4.3, yields

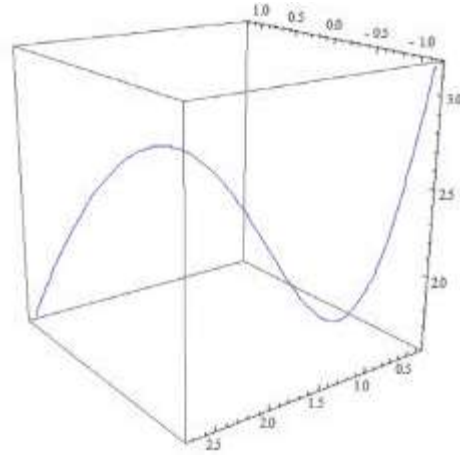


Fig. 1:

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